

## Calcium Signaling in Restricted Diffusion Spaces

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**ABSTRACT** One- and two-dimensional models of  $\text{Ca}^{2+}$  diffusion and regulation were developed and used to study the magnitudes and the spatial and temporal characteristics of the  $\text{Ca}^{2+}$  transients that are likely to develop in smooth muscle cells in restricted diffusion spaces between the plasma membrane and intracellular organelles. Simulations with the models showed that high  $[\text{Ca}^{2+}]$  (on the order of several  $\mu\text{M}$ ) can develop in such spaces and persist for 100–200 ms. These  $\text{Ca}^{2+}$  transients could: 1) facilitate the coupling of  $\text{Ca}^{2+}$  influx to intracellular  $\text{Ca}^{2+}$  release; 2) provide a mechanism for the regulation of stored  $\text{Ca}^{2+}$  that does not affect the contractile state of smooth muscle; 3) locally activate specific signal transduction pathways, before, or without activating other  $\text{Ca}^{2+}$  dependent pathways in the central cytoplasm of the cell. The latter possibility suggests that independent enzymatic processes in cells could be differentially regulated by the same intracellular second messenger.

### INTRODUCTION

In previous work with one- and two-dimensional models of  $\text{Ca}^{2+}$  diffusion and regulation in smooth muscle cells, Kargacin and Fay (1991) suggested that high  $\text{Ca}^{2+}$  concentrations can develop near the plasma membrane in smooth muscle cells as a result of  $\text{Ca}^{2+}$  influx through membrane channels. High  $\text{Ca}^{2+}$  concentrations could also develop within the cytoplasm at sites where  $\text{Ca}^{2+}$  is released from the sarcoplasmic reticulum (SR). During the decline of a  $\text{Ca}^{2+}$  transient, localized regions of low  $[\text{Ca}^{2+}]$  were predicted to develop in cells where  $\text{Ca}^{2+}$  is taken back up into the SR and at sites on the plasma membrane where  $\text{Ca}^{2+}$  is pumped out of the cell. The fact that inhomogeneities in  $[\text{Ca}^{2+}]$  are likely to be present in smooth muscle cells during transient  $\text{Ca}^{2+}$  signals suggests that the spatial characteristics of a  $\text{Ca}^{2+}$  signal are important determinants of the response of the cell to the signal. For example, an external stimulus that invokes a substantial extracellular influx of  $\text{Ca}^{2+}$  might activate enzymatic processes near the membrane that are not activated by an external stimulus that primarily invokes the release of  $\text{Ca}^{2+}$  from intracellular storage sites. Calcium influx into different regions of a cell could also influence cellular responses. Recent work by Bading et al. (1993) indicates that  $\text{Ca}^{2+}$  influx into neuronal cells, through different types of  $\text{Ca}^{2+}$  channels, results in the regulation of different signaling pathways. This finding can be explained if one assumes that each type of channel is localized to a specific region or regions of the plasma membrane and that the enzymes involved in each of the pathways are also confined spatially to sub-membrane regions near the channels.

Localized  $\text{Ca}^{2+}$  signaling would be especially important and prevalent in cells at locations where intracellular structures and organelles come into close contact with the surface

membrane. At such sites, the physical presence of the structures would impede the free movement of  $\text{Ca}^{2+}$  into the central cytoplasm of the cell and, as a consequence, the  $\text{Ca}^{2+}$  concentrations developed in these restricted diffusion spaces would be higher and persist for longer times than those developed in regions where free diffusion could occur. The possibility that local high concentrations of  $\text{Ca}^{2+}$  or other ions can develop and influence cellular function in smooth muscle and other types of cells in regions where  $\text{Ca}^{2+}$  influx or release is directed into a narrow cytoplasmic space has been proposed by a number of investigators. It has been suggested that, when depleted of  $\text{Ca}^{2+}$ , the superficial SR in smooth muscle takes up a significant fraction of the  $\text{Ca}^{2+}$  moving into the cell from the extracellular space and, thereby, reduces the amount of  $\text{Ca}^{2+}$  entering the cytoplasm (reviewed by van Breemen and Saida, 1989; Sturek et al., 1992). A number of laboratories (Benham and Bolton, 1986; Ohya et al., 1987; Hume and LeBlanc, 1989; Désilets et al., 1989; Stehno-Bittel and Sturek, 1992) have noted the spontaneous activation of  $\text{Ca}^{2+}$ -activated  $\text{K}^{+}$  currents in smooth muscle, which appears to be the result of the spontaneous release of  $\text{Ca}^{2+}$  from intracellular storage sites. Benham and Bolton (1986) and Stehno-Bittel and Sturek (1992) have suggested that this release occurs in regions where the SR comes into close apposition to the plasma membrane. Local  $\text{Ca}^{2+}$  signaling in restricted diffusion spaces has also been postulated to occur in cardiac muscle (see Lederer et al., 1990; Leblanc and Hume, 1990) and neurons (see Smith and Augustine, 1988). The evidence for this type of signaling has been largely indirect, however, and the magnitudes and the spatial and temporal properties of the ionic signals in such spaces have not been explored experimentally.

Smooth muscle cells are organized in a manner that could promote and exploit local  $\text{Ca}^{2+}$  signaling. The membrane of the SR in smooth muscle can come into close apposition to the plasma membrane (see Devine et al., 1972; Somlyo, 1980; Gabella, 1983; Somlyo and Franzini-Armstrong, 1985). Where this occurs, the SR membrane can run parallel to the plasma membrane for distances of 1  $\mu\text{m}$  or more (Devine et al., 1972; Gabella, 1983), and periodic structures

Received for publication 29 September 1993 and in final form 16 March 1994.

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0006-3495/94/07/262/11 \$2.00

that span the cytoplasmic space between the two membranes have been noted (Devine et al., 1972; Gabella, 1983; Somlyo and Franzini-Armstrong, 1985). The close contact and the presence of connecting structures between the SR and the surface membrane in smooth muscle cells suggests that these sites are specialized to promote the regulation of both cytoplasmic and stored  $\text{Ca}^{2+}$ . The spanning proteins could be directly involved in coupling excitation of the plasma membrane with release of stored intracellular  $\text{Ca}^{2+}$ , as are the foot proteins in cardiac and skeletal muscle (see Fleischer and Inui, 1989). Alternatively, the structures in smooth muscle might serve, primarily, to keep the membranes in close apposition. In this case, the  $\text{Ca}^{2+}$  channels and pumps on the plasma membrane could interact with  $\text{Ca}^{2+}$  release sites and pumps on the SR membrane as a result of the diffusion of  $\text{Ca}^{2+}$  or other substances through the narrow cytoplasmic space separating the two membranes.

To understand how local  $\text{Ca}^{2+}$  signals might develop in cells and influence specific signal transduction pathways, it is important to have theoretical knowledge of the likely magnitudes, time courses, and influences that can act on intracellular ionic signals, especially those that might arise in confined spaces. In the work reported here, the diffusion models described previously (Kargacin and Fay, 1991) were modified to explore specifically the properties of restricted diffusion spaces and their effects on  $\text{Ca}^{2+}$  signaling in smooth muscle cells. To provide a basic framework for understanding such signaling, the models were constructed without the imposition of any properties on the cell cytoplasm in the restricted spaces that were not similarly present in the rest of the cell.

Results of simulations with the models indicate that very high  $\text{Ca}^{2+}$  concentrations (on the order of several  $\mu\text{M}$ ) are likely to develop in restricted diffusion spaces in smooth muscle cells. The high concentrations could persist for 100–200 ms after transient  $\text{Ca}^{2+}$  influx and/or release from the SR. These high  $\text{Ca}^{2+}$  concentrations could function to rapidly couple  $\text{Ca}^{2+}$  influx to  $\text{Ca}^{2+}$ -induced  $\text{Ca}^{2+}$  release from the SR and, thereby, compensate in part for the relatively low diffusion coefficient for  $\text{Ca}^{2+}$  in cells (see Allbritton et al., 1992). Local  $\text{Ca}^{2+}$  movement into a restricted space in smooth muscle cells could also provide a means by which the  $\text{Ca}^{2+}$  content of the SR could be regulated independently of changes in bulk cytoplasmic  $[\text{Ca}^{2+}]$  and might be a means by which different enzymatic processes in the cell could be differentially controlled by the same second messenger. Use of the models to investigate the possibility of directly measuring the magnitude and time course of  $\text{Ca}^{2+}$  signals in restricted spaces suggested that such transients would be only poorly resolved by the fluorometric methods currently in common use.

## MATERIALS AND METHODS

The basic components of the one- and two-dimensional models for  $\text{Ca}^{2+}$  diffusion and regulation in smooth muscle cells were described previously (Kargacin and Fay, 1991). These models were based, as much as possible, on available experimental results and incorporated equations describing

$\text{Ca}^{2+}$  influx and extrusion through the plasma membrane,  $\text{Ca}^{2+}$  uptake and release from the SR, and intracellular  $\text{Ca}^{2+}$  buffering. Only a general outline of the components included in the models will be presented here (for additional details about the equations and parameters used, the reader is referred to Kargacin and Fay, 1991). Note: in Eqs. 1–7 below,  $\text{Ca}$  denotes  $[\text{Ca}^{2+}]_{\text{free}}$ .

The diffusion equations for radial ( $r$ ; Equation 1) and both radial and axial ( $r, z$ ; Eq. 2) diffusion in the model smooth muscle cells are

$$\frac{\partial \text{Ca}}{\partial t} = \frac{1}{r} \frac{\partial}{\partial r} \left( r D \frac{\partial \text{Ca}}{\partial r} \right) + F(\text{Ca}, t, r) \quad (1)$$

and

$$\frac{\partial \text{Ca}}{\partial t} = \frac{1}{r} \frac{\partial}{\partial r} \left( r D \frac{\partial \text{Ca}}{\partial r} \right) + \frac{\partial}{\partial z} \left( D \frac{\partial \text{Ca}}{\partial z} \right) + F(\text{Ca}, t, r, z), \quad (2)$$

where  $F(\text{Ca}, t, r)$  and  $F(\text{Ca}, t, r, z)$  include the concentration, time, and position-dependent  $\text{Ca}^{2+}$  regulatory processes in the cell as described below (see also Fig. 1 in Kargacin and Fay, 1991). To improve the spatial resolution of the models over those described previously (Kargacin and Fay, 1991), the model cell (3- $\mu\text{m}$  radius) was divided into 120 25-nm concentric annuli rather than the 30 (100 nm) annuli used previously. For the two-dimensional simulations, the model cell was divided into 30 (100 nm) or 60 (50 nm) annuli and 30–60 (100 nm) length elements. Movement of  $\text{Ca}^{2+}$  into and out of each spatial element in the models was computed using the explicit finite-difference formulation described by Crank (1975); see also Kargacin and Fay (1991). The smooth muscle plasma membrane  $\text{Ca}^{2+}$  pump was described by the Hill equation as

$$\frac{\Delta \text{Ca}}{\Delta t} = \frac{V_{\max}(\text{Ca})^n}{K_m^n + (\text{Ca})^n} \quad (3)$$

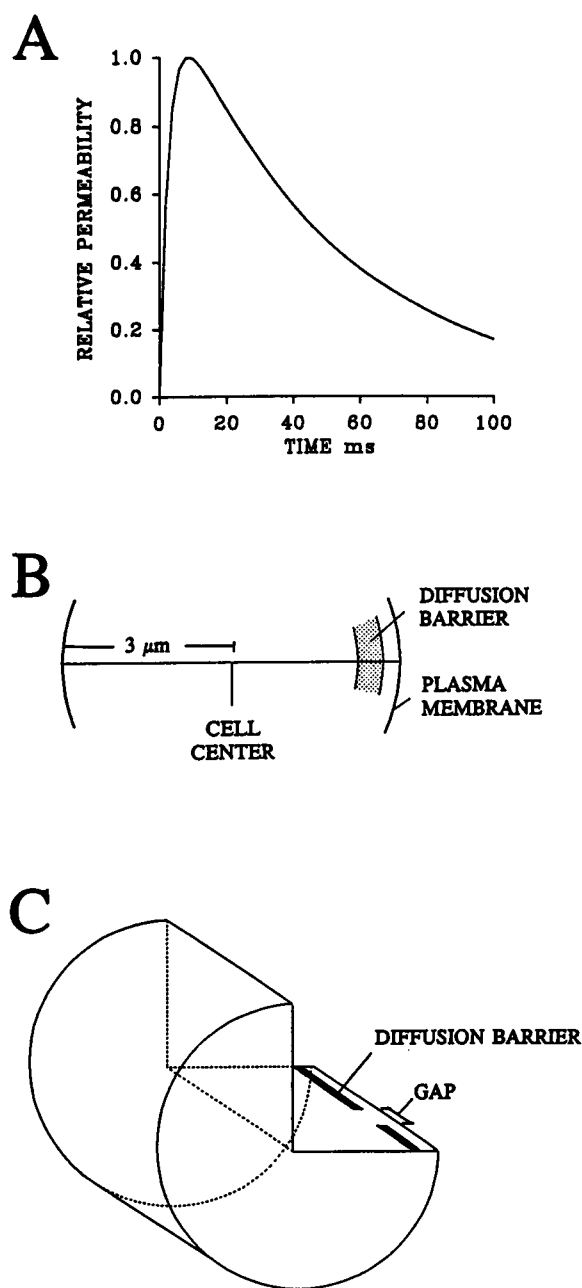
with  $V_{\max}$  the maximum velocity of the pump,  $K_m$  the  $[\text{Ca}^{2+}]$  at half maximal velocity, and  $n$  the Hill coefficient. To balance the pump efflux in the resting cell so that no net removal of  $\text{Ca}^{2+}$  from the cytoplasm occurred, an inward  $\text{Ca}^{2+}$  leak was included in the simulations and was described by the equation

$$\frac{\Delta \text{Ca}}{\Delta t} = K_{\text{leak}}(\text{Ca}_{\text{extracellular}} - \text{Ca}_{\text{cytoplasm}}) \quad (4)$$

with  $K_{\text{leak}}$  adjusted so that the  $\text{Ca}^{2+}$  influx through the leak was equal to the resting extrusion of  $\text{Ca}^{2+}$  through the plasma membrane. A similar pair of equations was used to describe the  $\text{Ca}^{2+}$  pump and a leak on the SR membrane when these elements were included in the model. In the latter case,  $\text{Ca}_{\text{extracellular}}$  in Eq. 4 was replaced by the  $[\text{Ca}^{2+}]$  in the SR (1.5 mM), and  $K_{\text{leak}}$  was adjusted to balance resting SR  $\text{Ca}^{2+}$  uptake. For the plasma membrane pump,  $V_{\max}$  was  $3.2 \times 10^{-13}$  mol/cm<sup>2</sup> s,  $K_m$  was 200 nM, and  $n$  was 1 (see Kargacin and Fay, 1991; Lucchesi et al., 1988; Carafoli, 1987). For the SR pump,  $V_{\max}$  was set at  $3.5 \times 10^{-12}$  mol/cm<sup>2</sup> s as described in Kargacin and Fay (1991). The other constants for the SR were  $K_m = 219$  nM and  $n = 2$  (see Kargacin et al., 1988). The extent to which  $\text{Na}^+/\text{Ca}^{2+}$  exchange contributes to  $\text{Ca}^{2+}$  regulation in smooth muscle cells is not well known (see Blaustein et al., 1991). To account for a possible contribution of the exchanger to  $\text{Ca}^{2+}$  extrusion, however, the velocity of the plasma membrane pump was doubled in the simulations with the model. This was based on the results of Cooney et al. (1991) (see also Kargacin and Fay, 1991), who suggest that  $\text{Na}^+/\text{Ca}^{2+}$  exchange contributes to  $\text{Ca}^{2+}$  extrusion to an extent equal to that of the plasma membrane pump. The  $\text{Ca}^{2+}$  buffers in the model cell were combined into a single fixed buffer distributed throughout the cell (total buffer concentration = 230  $\mu\text{M}$ ; see Bond et al., 1984; Allbritton et al., 1992). Buffering was described by the equation

$$\frac{\Delta \text{Ca}}{\Delta t} = -K_{\text{on}}([\text{buffer}]_{\text{free}})(\text{Ca}) + K_{\text{off}}[\text{Ca}]_{\text{bound}}, \quad (5)$$

with  $K_{\text{on}}$  and  $K_{\text{off}}$  the on and off rates of the buffer set at  $10^8/\text{M}\cdot\text{s}$  and  $10^2/\text{s}$ , respectively (see Robertson et al., 1981; Kargacin and Fay, 1991).  $\text{Ca}^{2+}$  influx into the cell through the plasma membrane and into the cytoplasmic space from intracellular  $\text{Ca}^{2+}$  stores were described by equations



**FIGURE 1** Time course of the  $\text{Ca}^{2+}$  permeability change and other elements incorporated into the one- and two-dimensional diffusion models. (A) Membrane permeability changes. Influx of  $\text{Ca}^{2+}$  through the plasma membrane and efflux from the SR were described by Eqs. 6 and 7. The time constants were adjusted as described in the text. (B) One-dimensional model. The cell diameter was 6  $\mu\text{m}$ ; extracellular  $[\text{Ca}^{2+}]$  was 1.5 mM; starting intracellular  $[\text{Ca}^{2+}]_{\text{free}}$  was 150 nM. The shaded area represents a barrier to free diffusion that was imposed between the central cytoplasm of the model cell and that near the plasma membrane. The thickness of the barrier was 100 nm in all simulations. The distance between the barrier and the plasma membrane was varied in different simulations. When the barrier was assumed to be due to the presence of the SR near the plasma membrane, SR  $\text{Ca}^{2+}$  pumps, leak, and release sites were located either on the cytoplasmic side or the plasma membrane side of the barrier. The  $[\text{Ca}^{2+}]_{\text{free}}$  inside the SR was assumed to be 1.5 mM. Components of the model describing  $\text{Ca}^{2+}$  influx, extrusion, and leakage through the plasma membrane and cytoplasmic  $\text{Ca}^{2+}$  buffering are discussed in the text. (C) Two-dimensional model. The  $\text{Ca}^{2+}$  regulatory components included in the two-dimensional model were the same as those used in the one-dimensional model, but both radial

of the form

$$\frac{\Delta \text{Ca}}{\Delta t} = K(\text{Ca}_{\text{out}} - \text{Ca}) \quad (6)$$

and

$$K = K_0(1 - \exp^{-t/t_{\text{on}}})(\exp^{-t/t_{\text{off}}}) \quad (7)$$

(see Cannell and Allen, 1984; Backx et al., 1989; Kargacin and Fay, 1991). The parameters  $t_{\text{on}}$  and  $t_{\text{off}}$  (3 and 50 ms, respectively) for plasma membrane influx were adjusted so that the time course of  $\text{Ca}^{2+}$  influx in the model cell matched the time course of the  $\text{Ca}^{2+}$  current measured by Becker et al. (1989) in voltage-clamped smooth muscle cells. The constant  $K_0$  was adjusted so that the average  $[\text{Ca}^{2+}]_{\text{free}}$  reached in the model cell (without diffusion barriers present) matched the  $[\text{Ca}^{2+}]_{\text{free}}$  (400–600 nM) measured experimentally in smooth muscle cells during  $\text{Ca}^{2+}$  transients. Adjustments to Eq. 7 for the two-dimensional model when small local currents (e.g., through single  $\text{Ca}^{2+}$  channels or clusters of channels) were modeled will be discussed later. The equation describing  $\text{Ca}^{2+}$  release from the SR had the same form and time constants as that used to describe influx through the plasma membrane except that, before release could start, the  $[\text{Ca}^{2+}]_{\text{free}}$  at the SR membrane had to reach a preset level ([switch]). The time course of the SR and plasma membrane permeability changes is shown in Fig. 1 A. The equations describing  $\text{Ca}^{2+}$  influx through the plasma membrane,  $\text{Ca}^{2+}$  buffering, extrusion out of the cell, and the inward leak of  $\text{Ca}^{2+}$  through the plasma membrane were included in all simulations. Uptake and release of  $\text{Ca}^{2+}$  by the SR was only included in simulations in which these processes were studied. The positions of the SR uptake and release sites in the model cell were varied in different simulations.

In addition to the elements described above, the models were constructed to include a diffusion barrier (a region with a lower diffusion coefficient) near the surface membrane of the cell (Fig. 1, B and C). Diffusion through the barrier region was hindered by the lower diffusion coefficient and, as a result, a restricted diffusion space was created between the surface membrane and the barrier. The boundaries on either side of the barrier were treated as described by Crank (1975) for diffusion through composite media. The thickness of the cytoplasmic space between the barrier and the surface membrane of the model cell was varied in different simulations. Although it was possible in the model to prevent completely the movement of  $\text{Ca}^{2+}$  into the central cytoplasmic space behind the barrier by setting the diffusion coefficient through the barrier region equal to 0, it was more realistic to allow some  $\text{Ca}^{2+}$  to reach the central cytoplasmic space to account for  $\text{Ca}^{2+}$  movement along pathways that circumvented the barrier itself (i.e., pathways that involved the length and angle dimensions of the cell that were not modeled). For an impenetrable barrier of thickness  $b$  and length  $2 \times l$ , using the equation  $r^2 = 2Dt$  to express the mean-square displacement of an ion diffusing through a medium with a diffusion coefficient  $D$  (see, for example, Hille, 1992), it can be seen that the amount of time ( $t$ ) required for  $\text{Ca}^{2+}$  to move around the barrier (from the center of one side to the center of the other side; path length =  $2l + b$ ) is equal to the amount of time required for  $\text{Ca}^{2+}$  to diffuse directly through the barrier region (path length =  $b$ ) with a diffusion coefficient of  $D_1 = Db^2/(2l + b)^2$ . For the simulations with the one-dimensional model described here, it was assumed that a typical intracellular barrier was 100 nm in thickness and extended parallel to the plasma membrane 100 nm in the length and angular directions beyond the site of  $\text{Ca}^{2+}$  influx. This assumption is consistent with the dimensions of the superficial SR in smooth muscle cells (see, for example, Gabella, 1983; Somlyo, 1980). To approximate this case, the diffusion coefficient through the barrier region in the model cell was set, therefore, at  $0.1 \times$  the coefficient in the rest of the cell.

The two-dimensional model (Fig. 1 C) was constructed with the same basic parameters as the one-dimensional model and permitted  $\text{Ca}^{2+}$

diffusion and diffusion along the cell length were modeled. The barrier to free diffusion (dark bars) in the two-dimensional model was not continuous in the length dimension but contained a gap (space between dark bars) through which free diffusion of  $\text{Ca}^{2+}$  into the central cytoplasm of the cell could occur.

movement to be studied simultaneously in both the radial direction and along the length of the cell. In the two-dimensional model, the diffusion barrier was not necessarily continuous along the cell length but could be interrupted by one or more gaps through which free diffusion into the bulk of the cytoplasm could occur. As was the case with the one-dimensional model, some diffusion through the barrier region was also allowed to occur to account for the movement of  $\text{Ca}^{2+}$  into the central cytoplasm through pathways involving the angular dimension. With the two-dimensional model, it was thus possible to look at restricted spaces in more global settings because movement through both restricted-diffusion regions and free-diffusion regions could be examined at the same time.

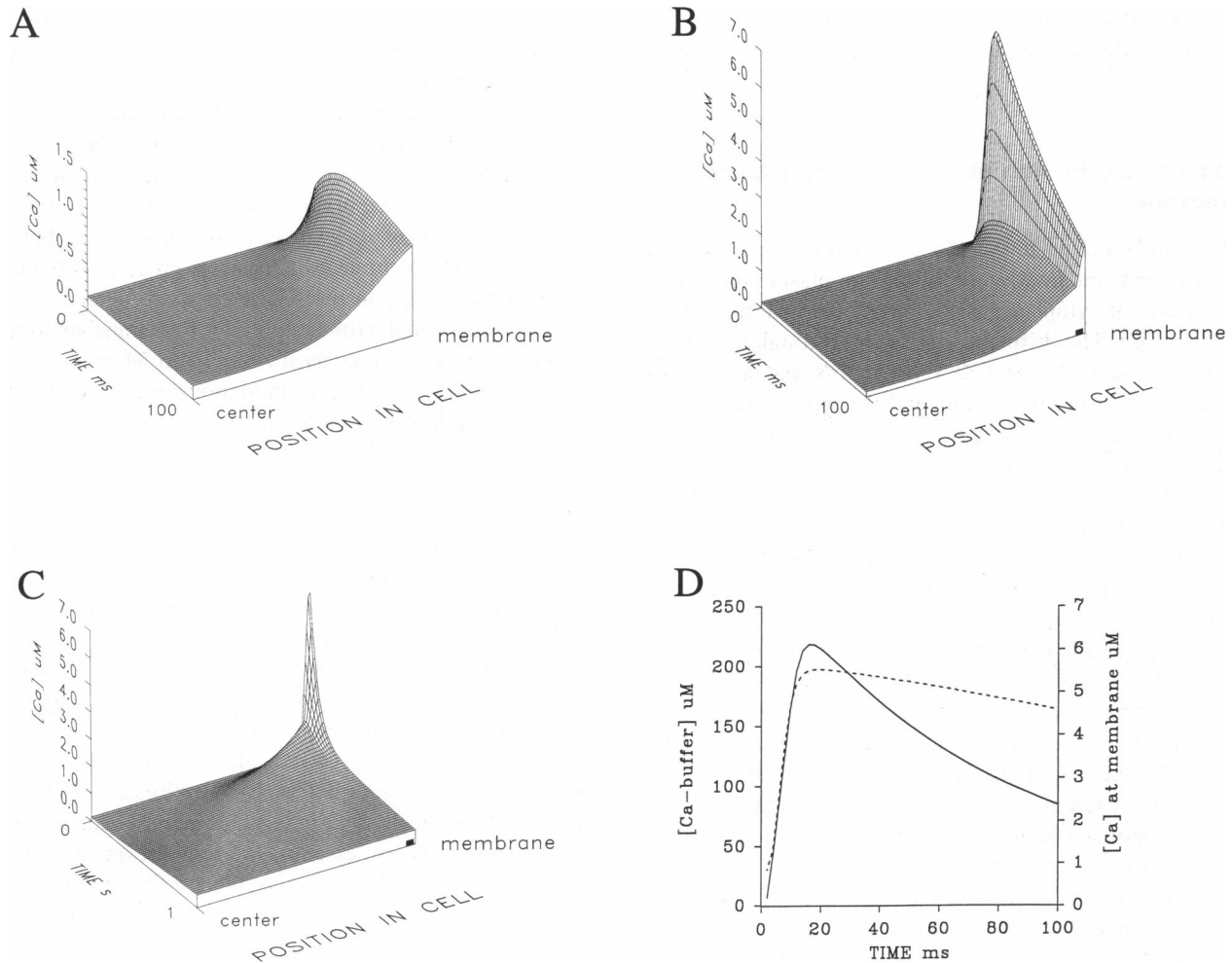
A Fortran program written by the author incorporating Eqs. 3–7 into the explicit finite differences formulae described by Crank (1975) (see also Kargacin and Fay, 1991) was used to obtain numerical solutions to Eqs. 1 and 2. The simulations were run on a 386 personal computer.

## RESULTS AND DISCUSSION

### Diffusion without a barrier

In previous work (Kargacin and Fay, 1991), an intracellular diffusion coefficient of  $4 \times 10^{-6} \text{ cm}^2/\text{s}$  (based on estimates

of the space not occupied by structures in electron micrographs of smooth muscle cells) was used to describe the movement of  $\text{Ca}^{2+}$  through the cytoplasm of the model cell. This number is approximately one-half the value of the diffusion coefficient for  $\text{Ca}^{2+}$  in water ( $7 \times 10^{-6} \text{ cm}^2/\text{s}$ ; see Kushmerick and Podolsky, 1969). In Fig. 2 A, the predicted  $[\text{Ca}^{2+}]_{\text{free}}$  profile in a smooth muscle cell without any barriers to diffusion and with an intracellular diffusion coefficient of  $4 \times 10^{-6} \text{ cm}^2/\text{s}$  is shown as a function of time. It can be seen that, even without an intracellular barrier, a substantial gradient of  $\text{Ca}^{2+}$  developed and persisted for more than 100 ms in the model cell after the start of  $\text{Ca}^{2+}$  influx through the surface membrane. A maximum  $[\text{Ca}^{2+}]_{\text{free}}$  of  $1.13 \mu\text{M}$  was reached just inside the plasma membrane in the cell 50 ms after the simulation started. During the 100 ms time period modeled, the  $[\text{Ca}^{2+}]_{\text{free}}$  in the cell halfway between the center of the cell and the plasma membrane rose from 150 to 212 nM.



**FIGURE 2** (A) Time course of radial diffusion without a barrier.  $[\text{Ca}^{2+}]_{\text{free}}$  profiles in the model cell are shown as a function of time with an intracellular diffusion coefficient of  $4 \times 10^{-6} \text{ cm}^2/\text{s}$  (A). (B–D) Radial diffusion with a barrier. The simulations were similar to those in A, but a 100-nm barrier to free diffusion located 25 nm from the plasma membrane was included in the model (the dark bar on the radial axis shows the location of the barrier relative to the plasma membrane). The diffusion coefficient through the barrier region was  $0.22 \times 10^{-6} \text{ cm}^2/\text{s}$ ; the coefficient in the cytoplasm on either side of the barrier was  $2.2 \times 10^{-6} \text{ cm}^2/\text{s}$ . The  $[\text{Ca}^{2+}]_{\text{free}}$  profiles in the model cell are shown at two time scales: (B) 100 ms total time, profiles plotted at 2-ms intervals; (C) 1 s total time, profiles plotted at 20-ms intervals. The average  $[\text{Ca}^{2+}]_{\text{free}}$  in the cell was 516 nM after 1 s. (D)  $[\text{Ca}^{2+}]_{\text{free}}$  versus time in the restricted space between the plasma membrane and the barrier (—; right scale) and  $\text{Ca}^{2+}$  bound to the  $\text{Ca}^{2+}$  buffer in this space versus time (---; left scale). For all of the simulations shown, the cell radius was  $3 \mu\text{m}$ , and the cell was divided into 120 concentric 25-nm annuli.

The recent work of Allbritton et al. (1992) indicates that the diffusion coefficient for  $\text{Ca}^{2+}$  in cytosolic extracts in which intracellular  $\text{Ca}^{2+}$  buffers are saturated is approximately  $2.2 \times 10^{-6} \text{ cm}^2/\text{s}$ ; roughly one-half the value used in the previous simulation ( $4 \times 10^{-6} \text{ cm}^2/\text{s}$ ). With a diffusion coefficient of  $2.2 \times 10^{-6} \text{ cm}^2/\text{s}$ , a maximum  $[\text{Ca}^{2+}]_{\text{free}}$  of  $1.55 \mu\text{M}$  was reached near the plasma membrane in the model cell 48 ms after  $\text{Ca}^{2+}$  influx began, and the  $[\text{Ca}^{2+}]_{\text{free}}$  halfway between the plasma membrane and the center of the cell rose from 150 to 167 nM during a 100 ms time period. Thus, as expected, the lower diffusion coefficient resulted in a steeper  $\text{Ca}^{2+}$  gradient near the plasma membrane and slower movement of  $\text{Ca}^{2+}$  into the interior of the cell. The simulations with both diffusion coefficients indicate, however, that the  $[\text{Ca}^{2+}]_{\text{free}}$  near the plasma membrane of a smooth muscle cell is likely to get quite high (1–2  $\mu\text{M}$ ) after extracellular  $\text{Ca}^{2+}$  influx even in the absence of structural barriers to free diffusion in the cytoplasm. The results further suggest that  $\text{Ca}^{2+}$ -dependent signal transduction processes are activated near the plasma membrane of a cell before other such processes are effected in the bulk of the cytoplasm.

In the simulations that follow, a cytoplasmic diffusion coefficient of  $2.2 \times 10^{-6} \text{ cm}^2/\text{s}$  was used.

### Effects of a diffusion barrier near the plasma membrane

In the simulations in Fig. 2 B–D, a 0.1  $\mu\text{m}$  thick barrier to free diffusion was imposed between the plasma membrane and the central cytoplasm of the model smooth muscle cell (see Fig. 1 B). The diffusion coefficient through the barrier region was assumed to be  $0.22 \times 10^{-6} \text{ cm}^2/\text{s}$ , and the distance between the plasma membrane and the barrier was varied. With the barrier 25 nm from the plasma membrane,  $[\text{Ca}^{2+}]_{\text{free}}$  in the restricted diffusion space between the plasma membrane and the barrier reached a maximum of 6.1  $\mu\text{M}$ , 16 ms after extracellular  $\text{Ca}^{2+}$  influx began. The  $[\text{Ca}^{2+}]_{\text{free}}$  in the restricted space then decreased as influx declined,  $\text{Ca}^{2+}$  binding to the cytoplasmic  $\text{Ca}^{2+}$  buffer occurred, and  $\text{Ca}^{2+}$  diffused into the central cytoplasm of the cell. Fig. 2 B shows the  $[\text{Ca}^{2+}]_{\text{free}}$  profile in the model cell during the first 100 ms after the start of  $\text{Ca}^{2+}$  influx, and Fig. 2 C shows the profile during the first 1 s. As can be seen in Fig. 2 C, although a substantial  $\text{Ca}^{2+}$  gradient developed in the restricted space under the conditions of the simulation, the very high  $[\text{Ca}^{2+}]_{\text{free}}$  was present for only 100–200 ms, roughly as long as influx continued (compare Figs. 1 A and 2 B). The rate at which the  $[\text{Ca}^{2+}]_{\text{free}}$  near the membrane declined after reaching a maximum value indicates that, once influx slowed, free  $\text{Ca}^{2+}$  was rapidly removed from the restricted space. There were three mechanisms in the model that could account for this: (1) extrusion of  $\text{Ca}^{2+}$  through the plasma membrane; (2) binding of  $\text{Ca}^{2+}$  to the  $\text{Ca}^{2+}$  buffer in the restricted space; (3) diffusion of  $\text{Ca}^{2+}$  through the barrier region into the central cytoplasm of the cell. The rate of  $\text{Ca}^{2+}$  extrusion (based on estimates of the maximum velocity of the plasma membrane  $\text{Ca}^{2+}$  pump and the likely contribution of  $\text{Na}^+/\text{Ca}^{2+}$  exchange; see Materials and Methods) was too slow to ac-

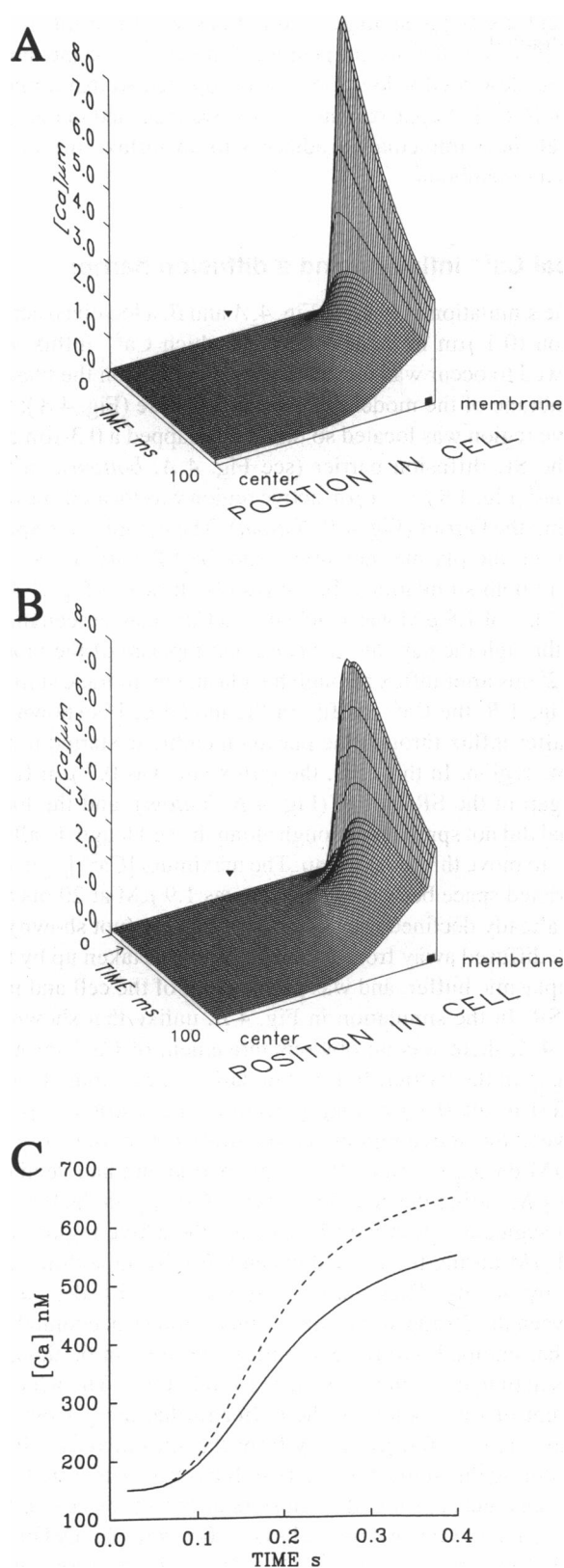
count for the rate of the decline. The  $\text{Ca}^{2+}$  transient near the membrane obtained when  $\text{Ca}^{2+}$  extrusion was not included in the model was virtually identical to the transient obtained when it was included. If the rate at which  $\text{Ca}^{2+}$  binding to the buffer occurred was slow relative to the rate of  $\text{Ca}^{2+}$  influx, the  $[\text{Ca}^{2+}]_{\text{free}}$  in the restricted space would decline after influx started to decrease as  $\text{Ca}^{2+}$  continued to bind to the buffer. Fig. 2 D compares the free  $\text{Ca}^{2+}$  transient in the restricted space for the simulation with the time course of  $\text{Ca}^{2+}$  binding to the buffer in the space. As can be seen, binding to the buffer was rapid enough to keep up with influx, and additional binding did not occur after the peak of the  $\text{Ca}^{2+}$  transient was reached. Thus, diffusion of  $\text{Ca}^{2+}$  out of the restricted space was the primary reason for the rapid decline of the  $\text{Ca}^{2+}$  transient. The relatively slow unloading of the buffer (Fig. 2 D) tended to keep  $\text{Ca}^{2+}$  in the restricted space. These results indicate that the diffusional properties of a restricted space and the extent of the space would be important determinants of how long a signal persisted in a space once  $\text{Ca}^{2+}$  influx declined. A low diffusion coefficient in this space or an extended barrier region would favor the confinement of  $\text{Ca}^{2+}$  in the space and the persistence of transient signals. The buffer in the space would retain  $\text{Ca}^{2+}$  and, as will be discussed later, might serve as a temporary  $\text{Ca}^{2+}$  store.

When the distance between the plasma membrane and the diffusion barrier was increased in the model, the maximum  $[\text{Ca}^{2+}]_{\text{free}}$  near the membrane decreased and the time required to reach this maximum increased. As the distance between the plasma membrane was increased from 25 to 100 nm, the maximum  $[\text{Ca}^{2+}]_{\text{free}}$  in the restricted space decreased from 6.1 to 4.9  $\mu\text{M}$  and the time required for the transient to reach its maximum value increased from 16 to 34 ms.

Although, in some of the simulations to follow, it will be assumed that the diffusion barrier was imposed by the physical presence of the SR near the plasma membrane, the work described to this point did not include SR  $\text{Ca}^{2+}$  uptake or release and, thus, the predicted high  $\text{Ca}^{2+}$  concentrations could develop anywhere in a cell where an intracellular structure is in close apposition to the plasma membrane. High  $\text{Ca}^{2+}$  concentrations such as those predicted by the model during  $\text{Ca}^{2+}$  transients could, in living smooth muscle cells, locally turn on  $\text{Ca}^{2+}$ -dependent processes near the membrane before the  $[\text{Ca}^{2+}]$  in the bulk of the cytoplasm changed significantly. The fact that the local  $\text{Ca}^{2+}$  signal did not persist in the model cell, however, indicates that, in considering this possibility, it is important to have estimates of the time courses of the local  $\text{Ca}^{2+}$  signals that might occur in smooth muscle cells. It is also essential to know the on and off rates and not just the equilibrium constants for  $\text{Ca}^{2+}$  binding to the intracellular effector proteins involved to determine whether such signaling is physiologically relevant.

### Coupling between $\text{Ca}^{2+}$ influx and SR $\text{Ca}^{2+}$ release in a restricted diffusion space

To examine the effects of diffusion barriers on the coupling between  $\text{Ca}^{2+}$  influx through the surface membrane and  $\text{Ca}^{2+}$  release from the SR,  $\text{Ca}^{2+}$  release sites were incorporated into



**FIGURE 3** Effect of a diffusion barrier on  $Ca^{2+}$ -induced  $Ca^{2+}$  release from the SR. (A, B)  $[Ca^{2+}]_{free}$  profiles versus time in the model cell with a 100-nm barrier to free diffusion located 25 nm from the plasma membrane (dark bar on radial axis as in Fig. 2). (A) SR  $Ca^{2+}$  release sites were located on the plasma membrane side of the barrier and released  $Ca^{2+}$  when the  $[Ca^{2+}]_{free}$  in the restricted space between the plasma membrane and the barrier reached  $1 \mu M$ . (B) SR  $Ca^{2+}$  release sites were located on the cytoplasmic side of the barrier. Release was triggered when the  $[Ca^{2+}]_{free}$  at

the model on either the cytoplasmic or the membrane side of the barrier. The cytoplasmic space between the plasma membrane and the SR was set at 25 nm. As in previous work (see Materials and Methods and Kargacin and Fay, 1991),  $Ca^{2+}$ -induced  $Ca^{2+}$  release from the SR was assumed to occur when the  $[Ca^{2+}]_{free}$  at the release sites reached a preset level ([switch]). Even with a [switch] of  $1 \mu M$  (see Iino 1989), when the release sites faced the plasma membrane,  $Ca^{2+}$  release from the SR was triggered very rapidly (release from the SR started 4 ms after influx began). In the results shown in Fig. 3 A, a maximum  $[Ca^{2+}]_{free}$  of  $8.3 \mu M$  was reached in the restricted space in the model cell 14 ms after influx began. A situation such as this in a smooth cell would facilitate a very rapid coupling of  $Ca^{2+}$  influx to SR release and might compensate somewhat for the low diffusion coefficient for buffered  $Ca^{2+}$  in the cell cytoplasm relative to the diffusion coefficients of other second messengers such as  $IP_3$  (see Allbritton et al., 1992). However, the  $Ca^{2+}$  released into a restricted diffusion space would be hindered from reaching the bulk of the cytoplasm by the physical presence of the barrier imposed by the SR itself. When the SR  $Ca^{2+}$  release sites in the model cell were located on the cytoplasmic side of the diffusion barrier, it took 20 ms for SR  $Ca^{2+}$  release to occur after the start of  $Ca^{2+}$  influx. The maximum  $[Ca^{2+}]_{free}$  that developed in the restricted space in Fig. 3 B was lower ( $6.1 \mu M$ ) than that seen in Fig. 3 A. In this case, SR release was not as rapidly coupled to influx, but  $Ca^{2+}$  moved into the central cytoplasm of the cell more rapidly than in the previous simulation. When  $Ca^{2+}$  was released into the restricted space, the  $[Ca^{2+}]_{free}$  in the central cytoplasm of the model cell (halfway between the cell center and the surface membrane) was 544 nM after 400 ms; when  $Ca^{2+}$  was released on the cytoplasmic side of the SR, the concentration at this same time point was 644 nM although, in the latter case, release of  $Ca^{2+}$  from the SR took 16 ms longer to start. Plots of  $[Ca^{2+}]_{free}$  versus time in the central cytoplasm of the model cell for the two simulations are shown in Fig. 3 C.

In previous work with the diffusion models (Kargacin and Fay, 1991), removal of  $Ca^{2+}$  from the cytoplasm by the SR  $Ca^{2+}$  pumps and extrusion of  $Ca^{2+}$  through the plasma membrane were predicted to have little influence on cellular  $[Ca^{2+}]$  during the early stages (the first 100 ms) of a transient  $Ca^{2+}$  signal. These processes occur on a longer time scale, and the gradients that developed in the model cell near intracellular  $Ca^{2+}$  uptake and plasma membrane extrusion sites were much smaller than those seen during  $Ca^{2+}$  influx and SR  $Ca^{2+}$  release and were apparent in the results only after  $Ca^{2+}$  had diffused throughout the cytoplasm. This was also true for the simulations described here. In simulations in

the sites reached  $1 \mu M$ . The SR release is apparent in B starting at approximately 20 ms (arrow on time axis). Other details are given in the text. (C)  $[Ca^{2+}]_{free}$  in the central cytoplasm of the model cell as a function of time for the two simulations. Cytoplasmic  $[Ca^{2+}]_{free}$  was monitored at a point midway between the cell membrane and the center of the cell (darts in A and B). Solid line in C: cytoplasmic  $[Ca^{2+}]_{free}$  versus time for the simulation in A. Dashed line in C: cytoplasmic  $[Ca^{2+}]_{free}$  versus time for the simulation in B.

which  $\text{Ca}^{2+}$ -induced  $\text{Ca}^{2+}$  release by the SR was allowed to occur into a restricted space (as in Fig. 3 A), the inclusion of an SR  $\text{Ca}^{2+}$  pump on the plasma membrane side of the SR did not noticeably alter the shape, magnitude, or time course of the  $\text{Ca}^{2+}$  transient in the restricted space over a time period of 100 ms. With the SR pump included in the model, a maximum  $[\text{Ca}^{2+}]_{\text{free}}$  of 8.32  $\mu\text{M}$  was reached in the restricted space after 14 ms; in the absence of the pump a maximum  $[\text{Ca}^{2+}]_{\text{free}}$  of 8.34  $\mu\text{M}$  was reached after 14 ms. There was also little difference in the SR  $\text{Ca}^{2+}$  content (assuming an initial SR  $[\text{Ca}^{2+}]_{\text{free}}$  of 1.5 mM; see Kargacin et al., 1988) because the  $\text{Ca}^{2+}$  release process in the SR essentially overwhelmed uptake during the early phases of the transient. It will be shown below that, when localized, low concentrations of  $\text{Ca}^{2+}$  were allowed to enter a restricted space, and the SR  $\text{Ca}^{2+}$  content could be increased without SR release.

### Simulations with the two-dimensional model

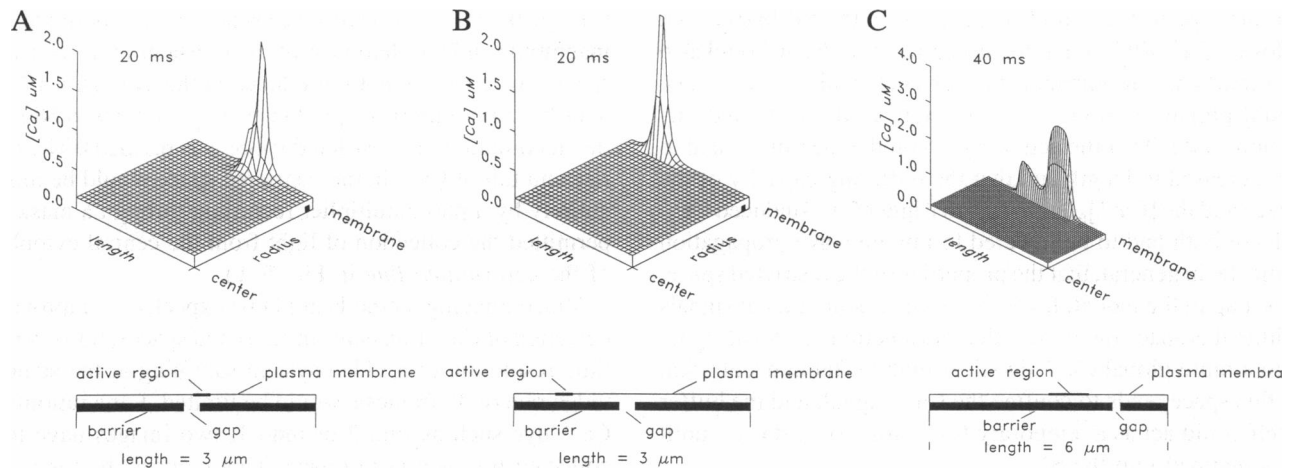
The simulations with the one-dimensional model described thus far provide information about the  $[\text{Ca}^{2+}]_{\text{free}}$  profile along the radius of a cylindrical cell that might be expected if movement of  $\text{Ca}^{2+}$  into the central cytoplasm of the cell was slowed by the presence of diffusion barriers. An assumption inherent in the model is that  $\text{Ca}^{2+}$  influx occurs over the entire cell surface and that the radial concentration profile seen in one cross section of the cell would be the same in other cross sections as one moved along the length of the cell. To study the  $\text{Ca}^{2+}$  signals that might arise as the result of localized  $\text{Ca}^{2+}$  influx into a smooth muscle cell, two-dimensional models were also developed. As described in Materials and Methods, the black bars in the two-dimensional model in Fig. 1 C represent a diffusion barrier, and the open space between the bars represents a gap in the barrier through which free diffusion of  $\text{Ca}^{2+}$  into the cytoplasm could occur. In the following simulations,  $\text{Ca}^{2+}$  influx through the plasma membrane was localized to different sites (active regions) along the membrane. The position of the active region, its size, and its position relative to the gap in the diffusion barrier were varied.

As discussed by Hille (1992), a density of open  $\text{Ca}^{2+}$  channels of approximately  $10/\mu\text{m}^2$  would be required to account for the macroscopic  $\text{Ca}^{2+}$  current measured in a typical cell assuming a physiological single channel current of 0.1 pA (see Hille, 1992; Gollasch et al., 1992). Thus, at the 100 nm spatial resolution of the two-dimensional model, assuming a uniform channel density of  $10/\mu\text{m}^2$ , one would expect at most a peak current of 0.1 pA passing through a single  $\text{Ca}^{2+}$  channel to enter a  $100\text{ nm} \times 100\text{ nm}$  spatial element (the possibility that  $\text{Ca}^{2+}$  channels are clustered at certain places in the plasma membrane will be discussed below). Stern (1992) modeled the  $\text{Ca}^{2+}$  concentration profile opposite a single channel assuming a 1 pA current through the channel. For such a channel, the Stern model predicted a  $[\text{Ca}^{2+}]$  of approximately 20  $\mu\text{M}$  at a position 50 nm (equivalent to the center of a length and radial element in the present model) from the plasma membrane. If one uses a simple linear scaling to scale down the prediction of Stern's model, a 0.1 pA

current under physiological conditions would result in a 2  $\mu\text{M}$   $[\text{Ca}^{2+}]$  at the 50 nm position. Therefore, for the simulations described below, Eq. 7 was adjusted so that a maximum  $[\text{Ca}^{2+}]$  of approximately 2  $\mu\text{M}$  was reached in the spatial element immediately adjacent to an influx site on the plasma membrane.

### Local $\text{Ca}^{2+}$ influx behind a diffusion barrier

In the simulations shown in Fig. 4, A and B, a localized active region (0.1  $\mu\text{m}$  in length) through which  $\text{Ca}^{2+}$  influx was allowed to occur was located at two positions on the plasma membrane of the model cell. In the first case (Fig. 4 A), the active region was located so that it overlapped a 0.3- $\mu\text{m}$  gap in the SR diffusion barrier (see Fig. 4 A, bottom); in the second (Fig. 4 B), a 0.1  $\mu\text{m}$  active region was located entirely behind the barrier (Fig. 4 B, bottom). The cytoplasmic space between the plasma membrane and the SR barrier was 0.1  $\mu\text{m}$  in both simulations. In the results shown in Fig. 4 A, a  $[\text{Ca}^{2+}]_{\text{free}}$  of 1.8  $\mu\text{M}$  was reached, and  $\text{Ca}^{2+}$  can be seen moving through the gap into the central cytoplasm of the model cell 20 ms after influx through the plasma membrane started. In Fig. 4 B, the  $\text{Ca}^{2+}$  profile in the model cell is shown 20 ms after influx through the plasma membrane started in the active region. In this case, the influx site was 0.7  $\mu\text{m}$  from the gap in the SR barrier (Fig. 4 B, bottom), and the local signal did not spread far enough along the cell length to allow  $\text{Ca}^{2+}$  to move through the gap. The maximum  $[\text{Ca}^{2+}]_{\text{free}}$  in the restricted space behind the barrier was 1.9  $\mu\text{M}$  at 20 ms but had already declined to 1.40  $\mu\text{M}$  after 40 ms (not shown) as  $\text{Ca}^{2+}$  diffused away from the influx site, was taken up by the cytoplasmic buffer, and was pumped out of the cell and into the SR. In the simulation in Fig. 4 B, unlike that shown in Fig. 4 A, there was no obvious movement of  $\text{Ca}^{2+}$  through the gap in the barrier. In this simulation, the average  $[\text{Ca}^{2+}]_{\text{free}}$  in the 0.4  $\mu\text{m}$  long portion of the restricted space, between the active region and the diffusion barrier, rose by 85  $\mu\text{M}$  during the first 20 ms time period and had declined by 7  $\mu\text{M}$  during the next 20 ms. The  $[\text{Ca}^{2+}]_{\text{free}}$  in the 0.4  $\mu\text{m}$  long segment within the SR opposite the active region rose by 4  $\mu\text{M}$  during the first 20 ms and rose by an additional 4  $\mu\text{M}$  by 40 ms. Thus, half of the  $\text{Ca}^{2+}$  lost by the buffer between the 20 and 40 ms time points could be accounted for by that pumped into the SR, and the buffer can be thought of as acting as a temporary store for this  $\text{Ca}^{2+}$ . The average amount of  $\text{Ca}^{2+}$  bound to the buffer further away from the influx site (0.2–0.6  $\mu\text{m}$  away from the site) increased by 3  $\mu\text{M}$  during this same time interval. It has been suggested that  $\text{Ca}^{2+}$  channels are found in clusters at certain places on the plasma membrane of some cells (see, for example, DeFelice, 1993; Risso and DeFelice, 1993). If this is the case, much higher local  $\text{Ca}^{2+}$  currents would occur at such locations. When the  $\text{Ca}^{2+}$  permeability of the plasma membrane in the model cell was increased by a factor of 10 over that used in the simulations described above (10-channel cluster), the  $[\text{Ca}^{2+}]_{\text{free}}$  in the restricted space for the simulation in Fig. 4 B was 18.1  $\mu\text{M}$  at 20 ms and 12.3  $\mu\text{M}$  at 40 ms. Note: a high  $[\text{Ca}^{2+}]_{\text{free}}$  such as this would also be seen at the SR membrane



**FIGURE 4**  $\text{Ca}^{2+}$  diffusion in two dimensions. (A, top) The  $[\text{Ca}^{2+}]_{\text{free}}$  in the model cell is shown 20 ms after a localized influx of  $\text{Ca}^{2+}$  through the plasma membrane began. (A, bottom) Schematic of the cell near the plasma membrane showing the relationship between the influx site (active region), an intracellular diffusion barrier (dark bars) and a gap in the barrier. In this case, the active region overlapped the location of the gap in the barrier and  $\text{Ca}^{2+}$  can be seen in the central cytoplasm of the cell in the region of the gap. (B, top) The  $[\text{Ca}^{2+}]_{\text{free}}$  in the model cell is shown 20 ms after  $\text{Ca}^{2+}$  influx through the active region started. In this case, the active region of the membrane was located entirely behind the diffusion barrier (B, bottom) and  $\text{Ca}^{2+}$  did not reach the gap or move into the central cytoplasm. For the simulations in A and B, the cell was divided into 30 100-nm concentric annuli and 30 100-nm length elements. The diffusion barrier was 100 nm thick and was located 100 nm from the plasma membrane. (C) Generation of a  $\text{Ca}^{2+}$  wave in the restricted space between the SR and the plasma membrane. For this simulation, the active region was 1  $\mu\text{m}$  long and was located entirely behind the diffusion barrier (C, bottom). The cell was divided into 30 100-nm concentric annuli and 60 100-nm length elements. Other details of the simulation are the same as those for the simulation in A and B or are described in the text. The  $[\text{Ca}^{2+}]_{\text{free}}$  profile in the cell is shown 40 ms after  $\text{Ca}^{2+}$  influx through the active region started. Note the  $\text{Ca}^{2+}$  wave on the left side of the active region propagating toward the upper left of the figure. For the simulations shown in A–C, the diffusion coefficient through the barrier region was  $0.22 \times 10^{-6} \text{ cm}^2/\text{s}$ ; the coefficient in the cytoplasm on either side of the barrier and through the gap was  $2.2 \times 10^{-6} \text{ cm}^2/\text{s}$ .

near a single  $\text{Ca}^{2+}$  channel if the distance between the SR and the plasma membrane was less than the 0.1  $\mu\text{m}$  separation that was used in the present simulations. In the 10-channel simulation, the average  $[\text{Ca}^{2+}]$  in the 0.4  $\mu\text{m}$  length of SR opposite the influx site increased by 4  $\mu\text{M}$ , and the average  $[\text{Ca}^{2+}\text{-buffer}]$  decreased from 210 to 202  $\mu\text{M}$  during the time interval between 20 and 40 ms. Although the higher  $\text{Ca}^{2+}$  influx resulted in a higher local  $[\text{Ca}^{2+}]_{\text{free}}$ , the  $\text{Ca}^{2+}$  signal in the model cell remained confined (qualitatively similar to the result shown in Fig. 4 B), and no obvious movement of  $\text{Ca}^{2+}$  through the gap in the barrier was visible.

Small  $\text{Ca}^{2+}$  influxes, such as those described above, into a restricted diffusion space between the plasma membrane and the SR in a smooth muscle cell could be a means of loading the SR. This could occur without a corresponding change in  $[\text{Ca}^{2+}]$  in the central cytoplasm large enough to elicit a contractile response in a cell. Conversely, small amounts of  $\text{Ca}^{2+}$  released by the SR into a restricted space could be pumped out of a cell as a means of regulating SR  $\text{Ca}^{2+}$  content without eliciting changes in the central cytoplasmic  $[\text{Ca}^{2+}]$  or contractile responses. These simulations suggest that localized  $\text{Ca}^{2+}$  influx behind a diffusion barrier could remain localized and could be used by the cell to regulate enzymes confined to such regions.

### Local $\text{Ca}^{2+}$ influx, SR $\text{Ca}^{2+}$ release, and the propagation of $\text{Ca}^{2+}$ waves

In the simulations with the two-dimensional model described above, the SR was allowed to take up  $\text{Ca}^{2+}$  from the restricted space opposite the active region of the plasma membrane, but

SR release was not included in the model. If the  $\text{Ca}^{2+}$  in the restricted space did trigger SR release, one might expect that the release could become regenerative and initiate a wave of  $\text{Ca}^{2+}$  release that propagated along the surface membrane in the space. In this way, a confined  $\text{Ca}^{2+}$  signal, due to a localized influx of  $\text{Ca}^{2+}$  large enough to trigger SR  $\text{Ca}^{2+}$  release, might lead to a more global intracellular release of  $\text{Ca}^{2+}$  and possibly a contractile response. To test this possibility, the diffusion barrier in the two-dimensional model was constructed as in the previous simulations shown in Fig. 4 B. In the present case, however, SR release was allowed to occur when the  $[\text{Ca}^{2+}]_{\text{free}}$  in the space reached 1  $\mu\text{M}$ . Although SR  $\text{Ca}^{2+}$  release into the restricted space was triggered by the  $\text{Ca}^{2+}$  influx expected from a single  $\text{Ca}^{2+}$  channel, even when a cluster of 10 channels was incorporated into the model, the  $\text{Ca}^{2+}$  signal did not propagate along the membrane (not shown but qualitatively similar to the  $\text{Ca}^{2+}$  signal in Fig. 4 B). When the total buffer concentration in the space was reduced from 230 to 115  $\mu\text{M}$ , the active region of the membrane was expanded to 1  $\mu\text{m}$  and the [switch] was lowered to 400 nM; the local  $\text{Ca}^{2+}$  influx did lead to the formation of a  $\text{Ca}^{2+}$  pulse or wave that propagated along the cell length in the restricted space in one direction (toward the upper left in Fig. 4 C). However a 0.2- $\mu\text{m}$  gap in the SR barrier (located 0.4  $\mu\text{m}$  away from the influx site; see schematic in Fig. 4 C) effectively prevented a wave from moving in the other direction (toward the lower right) in the cell. These results indicate that  $\text{Ca}^{2+}$  waves dependent on  $\text{Ca}^{2+}$ -induced  $\text{Ca}^{2+}$  release could propagate along the cell length through a restricted space but also indicate that the size of the space and its  $\text{Ca}^{2+}$  buffering capacity exert critical influences

on this type of  $\text{Ca}^{2+}$  movement. Gaps in the SR barrier that allowed  $\text{Ca}^{2+}$  diffusion into the central cytoplasm would also be important determinants of the extent to which a  $\text{Ca}^{2+}$  wave could propagate along the cell length in the restricted diffusion space. That the active region on the membrane had to be increased in length and that the buffering capacity of the space and the  $[\text{Ca}^{2+}]_{\text{free}}$  required to trigger  $\text{Ca}^{2+}$ -induced  $\text{Ca}^{2+}$  release both had to be lowered to initiate wave propagation indicate, in general, that the properties of the restricted space, as set up in the model, favor the confinement of local signals within this space rather than the transduction of a local signal into a more global  $\text{Ca}^{2+}$  signal. A high buffer concentration in this space tends to confine the  $\text{Ca}^{2+}$  signal, and the buffer itself could act as a temporary  $\text{Ca}^{2+}$  store to hold  $\text{Ca}^{2+}$  until it is pumped into the SR.

### Predictions for the measurement of intracellular $\text{Ca}^{2+}$ in experiments on smooth muscle

At present, many laboratories are actively involved in using fluorescent indicators to measure  $\text{Ca}^{2+}$  transients in smooth muscle cells. Fluorescence signals are either collected with a photomultiplier from a masked region of a cell or video images of the cell are acquired. To determine how a  $\text{Ca}^{2+}$  signal in a restricted space might be perceived with a non-imaging method of measurement, the average  $[\text{Ca}^{2+}]_{\text{free}}$  over the entire model cell for the simulation shown in Fig. 2 B–D and the average  $[\text{Ca}^{2+}]_{\text{free}}$  in the central cytoplasm of the cell (averaged over the central 4.5  $\mu\text{m}$  of a 6  $\mu\text{m}$  diameter cell) were compared with the  $\text{Ca}^{2+}$  transient in the restricted space. A photomultiplier recording through a mask that allowed the collection of light from the entire cell, including the area near the plasma membrane, would be expected to produce a signal similar to the average  $\text{Ca}^{2+}$  curve (*dashed line*) in Fig. 5 A. An initial transient peak in  $[\text{Ca}^{2+}]_{\text{free}}$  corresponding to the

$\text{Ca}^{2+}$  peak in the restricted space is present in this signal. The magnitude and the sharpness of the transient in the restricted space, however, are not reproduced in the average  $\text{Ca}^{2+}$  signal (the maximum  $[\text{Ca}^{2+}]$  in the average curve was 680 nM; the maximum in the restricted space was 6.1  $\mu\text{M}$ ). The rapid rise and fall of  $\text{Ca}^{2+}$  in the restricted space would be missed entirely by a photomultiplier recording through a mask that permitted the collection of light from the central cytoplasm of the cell (*dotted line* in Fig. 5 A).

Video imaging methods might be expected to improve the detection of  $\text{Ca}^{2+}$  transients in restricted spaces; however, the time resolution of a video system sampling at the standard video rate of 30 frames/s would be limited. For a ratiometric  $\text{Ca}^{2+}$  dye such as fura-2 or indo-1, two images have to be collected for each ratio image. This reduces the time resolution to 15 frames/s. Fig. 5 B (*solid line*) shows the  $\text{Ca}^{2+}$  transient in the restricted space (from Fig. 2 B) and the transient in the space (*filled circles* and *dashed line*) that would be recorded by a system (with sufficient spatial resolution to record  $\text{Ca}^{2+}$  in the 25 nm restricted space) that collected light and computed the  $[\text{Ca}^{2+}]_{\text{free}}$  in the space 17 times/s (the concentrations shown by the solid circles in Fig. 5 B are the average concentrations in the restricted diffusion space over each sampling period).

In the above discussion, it was assumed that a  $\text{Ca}^{2+}$  dye would not add additional  $\text{Ca}^{2+}$  buffering capacity to the model cell (for discussions of the effects of the addition of  $\text{Ca}^{2+}$  buffers to cells, see Stern, 1992; Blumenfeld et al., 1992). It was also assumed that light could be collected from a single image plane (high  $z$  axis resolution). Addition of  $\text{Ca}^{2+}$  buffer to a living smooth muscle cell would reduce the magnitude of and slow down a transient  $\text{Ca}^{2+}$  signal. Out of focus light from the area near the membrane of a cell above and below a central plane of focus might actually allow a  $\text{Ca}^{2+}$  transient near the membrane to be recorded from a

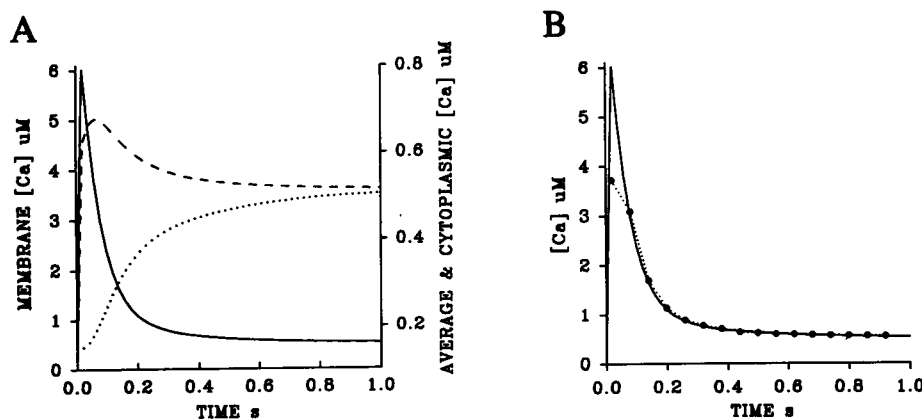


FIGURE 5 Simulated measurements of  $[\text{Ca}^{2+}]_{\text{free}}$ . (A)  $[\text{Ca}^{2+}]_{\text{free}}$  in the restricted space (—, left scale); average  $[\text{Ca}^{2+}]_{\text{free}}$  over the entire model cell (---, right scale); and average central cytoplasmic  $[\text{Ca}^{2+}]_{\text{free}}$  (···, right scale) for the simulation in Fig. 2 B. To calculate average  $[\text{Ca}^{2+}]_{\text{free}}$  over the entire model cell, the total free  $\text{Ca}^{2+}$  in each annulus in the model cell was computed and used to determine the total free  $\text{Ca}^{2+}$  in the cell. This value was then divided by the cell volume. To calculate the average central cytoplasmic  $[\text{Ca}^{2+}]_{\text{free}}$ , the total free  $\text{Ca}^{2+}$  in each annulus for the first 90 annuli in the cell (central 4.5  $\mu\text{m}$  of the 6  $\mu\text{m}$  diameter cell) was computed and used to determine the total free  $\text{Ca}^{2+}$  in this region. Total free  $\text{Ca}^{2+}$  was then divided by the volume of the region. Note: on the right hand scale in A average  $[\text{Ca}^{2+}] = \text{average } [\text{Ca}^{2+}]_{\text{free}}$  across the entire cell; cytoplasmic  $[\text{Ca}^{2+}] = \text{average } [\text{Ca}^{2+}]_{\text{free}}$  across the central cytoplasm of the model cell. (B)  $[\text{Ca}^{2+}]_{\text{free}}$  in the restricted space (—) and the  $[\text{Ca}^{2+}]_{\text{free}}$  (···●···) that would be reported by a measurement in which the  $[\text{Ca}^{2+}]_{\text{free}}$  in the restricted space was averaged over 59 ms time intervals and reported at a rate of 17/s.

measurement of  $\text{Ca}^{2+}$  in the central cytoplasm of a cell. The results presented above, however, indicate that, in general, such transients would be only poorly resolved experimentally. Until improvements in the temporal and spatial resolution of current experimental techniques for the direct measurement of intracellular  $\text{Ca}^{2+}$  are made, the most sensitive methods for detecting local  $\text{Ca}^{2+}$  changes near the plasma membrane in cells are probably those that indirectly do this by monitoring  $\text{Ca}^{2+}$ -dependent changes in ion channel activity.

## SUMMARY

The purpose of this study was to provide a general description of the temporal and spatial characteristics of the  $\text{Ca}^{2+}$  concentration gradients that might develop in smooth muscle cells after  $\text{Ca}^{2+}$  influx through the plasma membrane into regions where there are barriers that inhibit the free movement of  $\text{Ca}^{2+}$  into the central cytoplasm. The parameters describing the  $\text{Ca}^{2+}$  regulatory mechanisms that were built into the model were based as much as possible on those derived experimentally and were assumed to be the same throughout the cell. Results of the simulations indicate the following.

- Very high  $\text{Ca}^{2+}$  concentrations and steep gradients can develop in cells during transient  $\text{Ca}^{2+}$  signals, especially where free diffusion into the cytoplasm is restricted by the presence of physical barriers.
- The  $\text{Ca}^{2+}$  gradients and concentrations that accompany the influx of  $\text{Ca}^{2+}$  into the cell from the extracellular space might themselves be important elements of a  $\text{Ca}^{2+}$  signal and serve to couple influx to intracellular  $\text{Ca}^{2+}$  release. These local signals might also activate  $\text{Ca}^{2+}$ -dependent signal transduction pathways near the plasma membrane before triggering cell contraction. Local  $\text{Ca}^{2+}$  influx or release into a restricted space might also occur and raise the  $[\text{Ca}^{2+}]$  in the space without significantly changing the  $[\text{Ca}^{2+}]$  of the central cytoplasm of the cell.
- These local  $\text{Ca}^{2+}$  signals might serve to regulate intracellular  $\text{Ca}^{2+}$  stores or might provide a means by which the same second messenger (e.g.,  $\text{Ca}^{2+}$ ) could activate different signal transduction pathways in response to different stimuli.

The results as a whole indicate that it is important to consider intracellular second messenger signals not only in terms of their temporal characteristics but also in terms of how these signals might be spatially distributed in a cell.

The author wishes to thank Drs. Gisele Scott-Woo, Michael P. Walsh, and Margaret E. Kargacin for their critical comments on the manuscript. The author is an Alberta Heritage Foundation Scholar. This work was supported by National Institutes of Health grant AR39678 and the Heart and Stroke Foundation of Alberta.

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